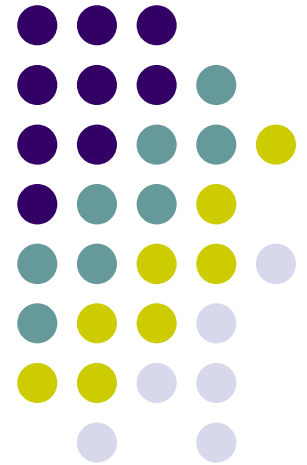


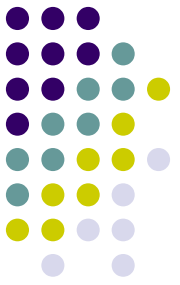
# Chapter 30

## Sources of the Magnetic Field



# PHYS 2321

## Week 10: Sources of Magnetic field



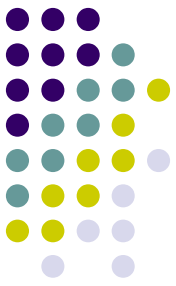
### Day 1 Outline

- 1) Hwk: Ch. 28 P. 1,4,5,19,27,29,37,38 Due Wed → Fri  
MiscQ 1-13 (odd)
- 2) Quiz on Ch. 27 – magnetic fields and forces
- 3) Applications of crossed E and B fields
- 4) Sources of magnetic fields (Ch 28)
  - a. Demo: current makes B-field
  - b. Biot-Savart law
  - c. B of a long straight wire
  - d. Force between parallel currents

Notes:

# PHYS 2321

## Week 10: Sources of Magnetic field



### Day 2 Outline

1) Hwk: Ch. 28 P. 1,4,5,19,27,29,37,38 Due Fri  
MiscQ 1-13 (odd)

2) Sources of magnetic fields (Ch 28)

a. Demo: current makes B-field ✓

b. Biot-Savart law

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{Id\vec{s} \times \hat{r}}{r^2}$$

c. B of a long straight wire

d. Force between parallel currents

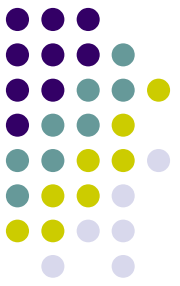
e. Ampere's law – easier than Biot-Savart!

Notes:

# Biot-Savart Law – Introduction

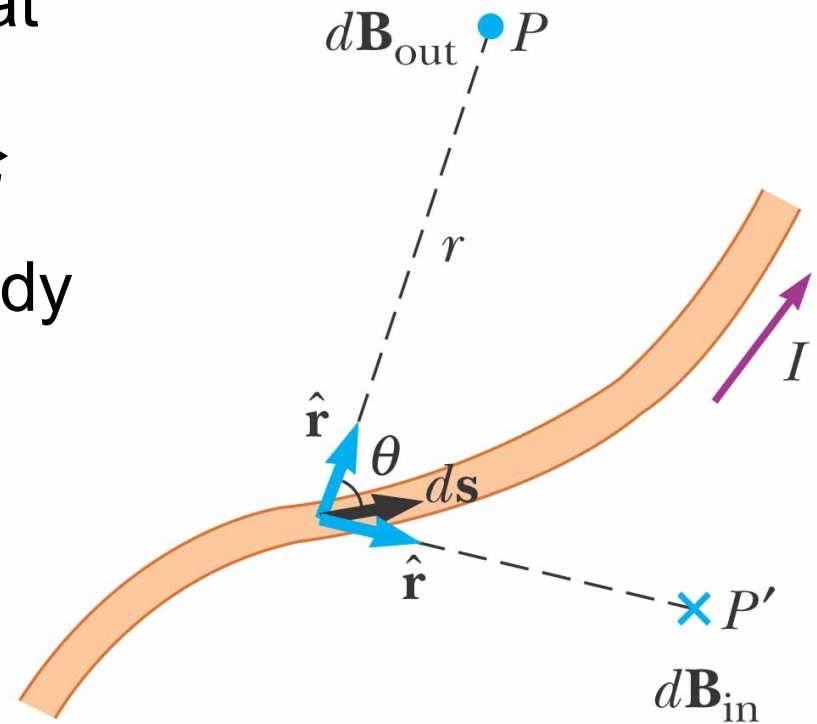


- Biot and Savart conducted experiments on the force exerted by an electric current on a nearby magnet
- They arrived at a mathematical expression that gives the magnetic field at some point in space due to a current

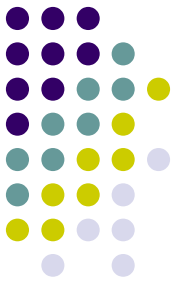


# Biot-Savart Law – Set-Up

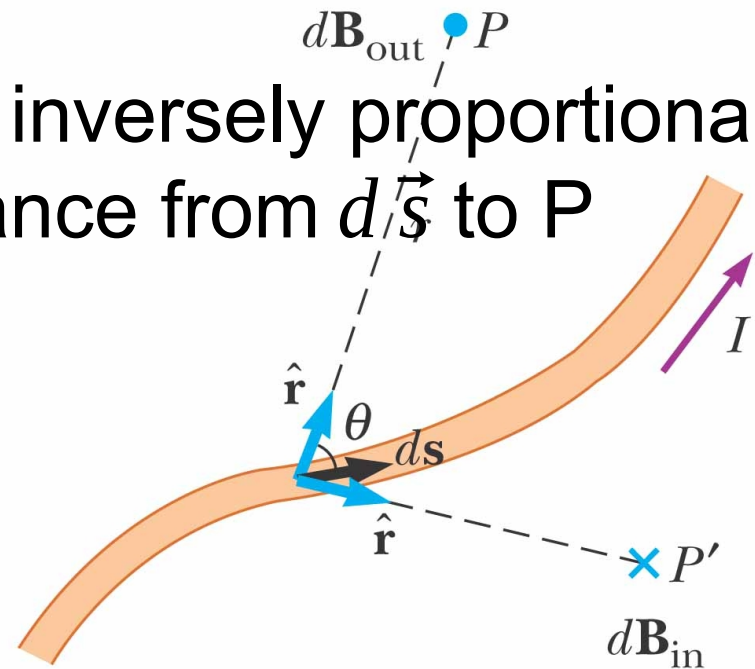
- The magnetic field is  $d\vec{B}$  at some point  $P$ .
- The length element is  $d\vec{s}$
- The wire is carrying a steady current,  $I$



# Biot-Savart Law – Observations



- The vector  $d\vec{B}$  is perpendicular to both  $d\vec{s}$  and to the unit vector  $\hat{r}$  directed from  $d\vec{s}$  toward P
- The magnitude of  $d\vec{B}$  is inversely proportional to  $r^2$ , where  $r$  is the distance from  $d\vec{s}$  to P



# Biot-Savart Law – Observations, cont



- The magnitude of  $d\vec{B}$  is proportional to the current and to the magnitude  $ds$  of the length element  $d\vec{s}$
- The magnitude of  $d\vec{B}$  is proportional to  $\sin \theta$ , where  $\theta$  is the angle between the vectors  $d\vec{s}$  and  $\hat{r}$



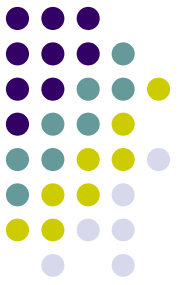
# Biot-Savart Law – Equation

- The observations are summarized in the mathematical equation called the **Biot-Savart law**:

$$d \vec{B} = \frac{\mu_0}{4 \pi} \frac{I d \vec{s} \times \hat{r}}{r^2}$$

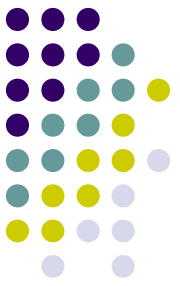
- The magnetic field described by the law is the field *due to* the current-carrying segment  $d \vec{s}$ .
  - It doesn't include the B due to other currents or permanent magnets.





# Permeability of Free Space

- The constant  $\mu_0$  is called the **permeability of free space**
- $\mu_0 = 4\pi \times 10^{-7} \text{ T} \cdot \text{m} / \text{A}$



# Total Magnetic Field

- $d\vec{B}$  is the field created by the current in the length segment  $ds$
- To find the total field, sum up the contributions from all the current elements  $I d\vec{s}$

$$\vec{B}_{tot} = \frac{\mu_0 I}{4\pi} \int_{current} \frac{d\vec{s} \times \hat{r}}{r^2}$$

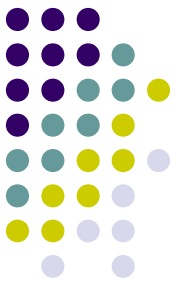
- The integral is over the entire current distribution

# $\vec{B}$ Compared to $\vec{E}$



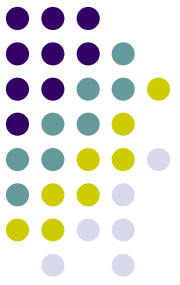
- Distance
  - The magnitude of the magnetic field varies as the inverse square of the distance from the source
  - The electric field due to a point charge also varies as the inverse square of the distance from the charge

# $\vec{B}$ Compared to $\vec{E}$ , 2



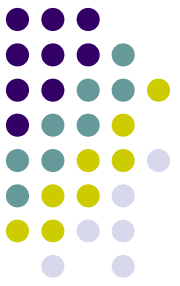
- Direction
  - The electric field created by a point charge is radial in direction
  - The magnetic field created by a current element is perpendicular to both the length element  $d\vec{s}$  and the unit vector  $\hat{r}$

# $\vec{B}$ Compared to $\vec{E}$ , 3



- Source
  - An electric field is established by an isolated electric charge
  - The current element that produces a magnetic field must be part of an extended current distribution
    - Therefore you must integrate over the entire current distribution

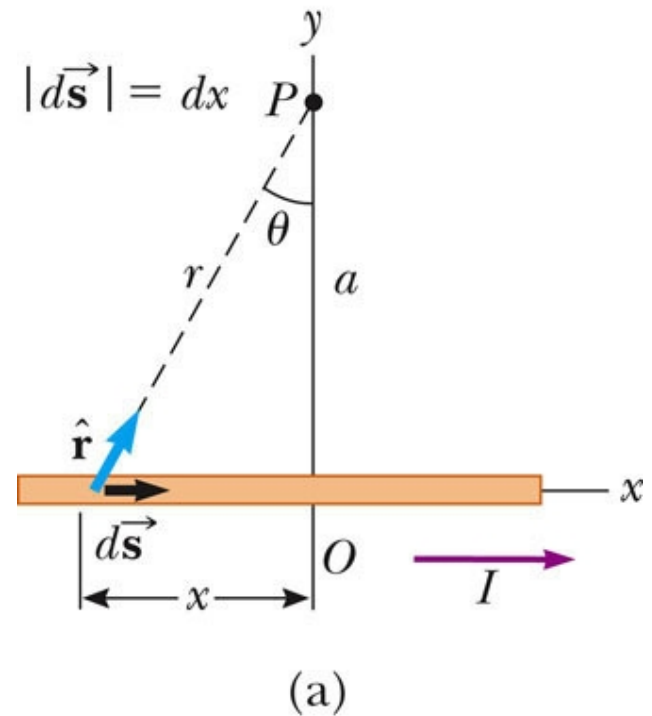
# $\vec{B}$ for a Long, Straight Conductor



- The thin, straight wire is carrying a constant current,  $I$
- $d\vec{s} \times \hat{r} = (dx \sin(90 - \theta)) \hat{k}$
- Integrating over all the current elements gives

$$B = -\frac{\mu_0 I}{4\pi a} \int_{\theta_1}^{\theta_2} \cos \theta d\theta$$

$$B = \frac{\mu_0 I}{4\pi a} (\sin \theta_1 - \sin \theta_2)$$

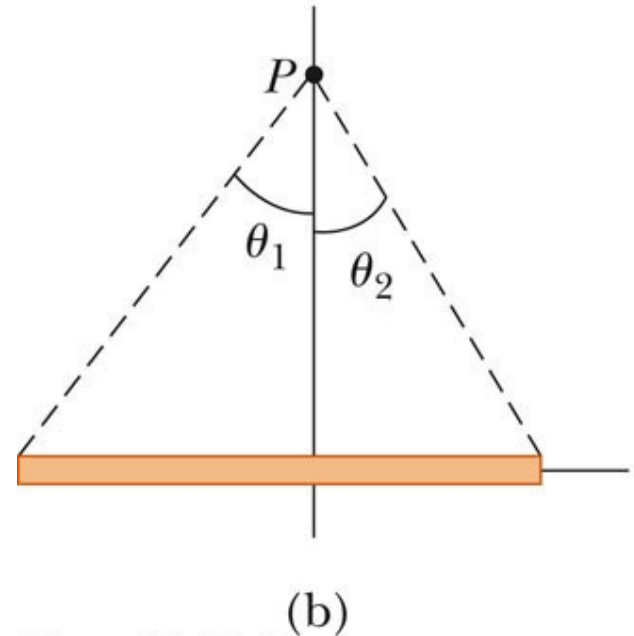


# $\vec{B}$ for a Long, Straight Conductor, Special Case

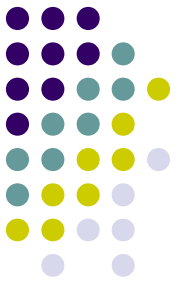


- If the conductor is an infinitely long, straight wire,  $\theta_1 = \pi/2$  and  $\theta_2 = -\pi/2$
- The field becomes

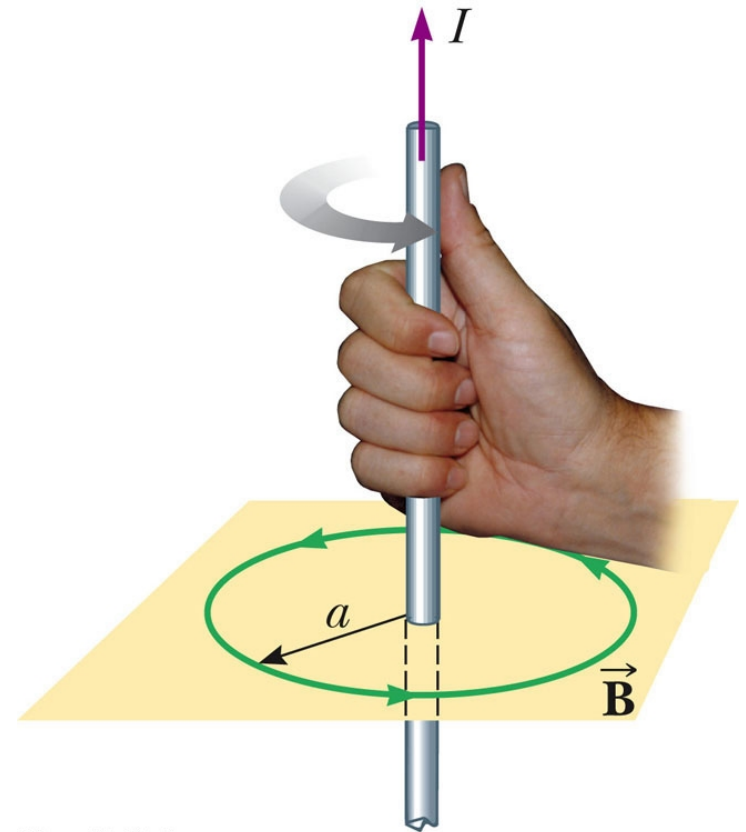
$$B = \frac{\mu_0 I}{2 \pi r}$$



# B for a Long, Straight Conductor, Direction

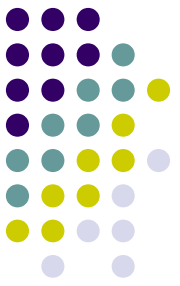


- The magnetic field lines are circles concentric with the wire
- The field lines lie in planes perpendicular to to wire
- The right-hand rule for determining the direction of the field is shown





# B for a Curved Wire Segment

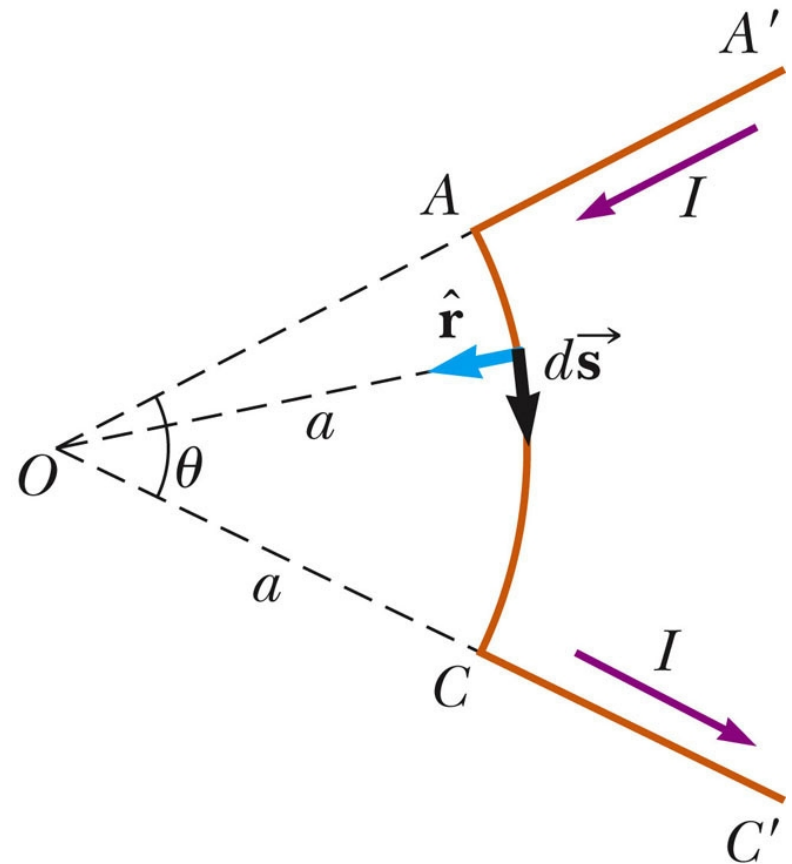


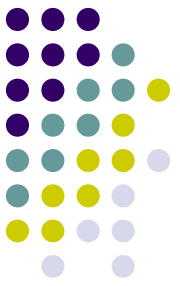
- Find the field at point  $O$  due to the wire segment

- $I$  and  $R$  are constants

$$B = \mu_0 I \theta / 4\pi a$$

$\theta$  will be in radians





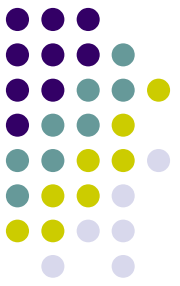
# B for a Circular Loop of Wire

- Consider the previous result, with a full circle

- $\theta = 2\pi$

$$B = \frac{\mu_0 I \theta}{4 \pi a} = \frac{\mu_0 I 2 \pi}{4 \pi a} = \frac{\mu_0 I}{2 a}$$

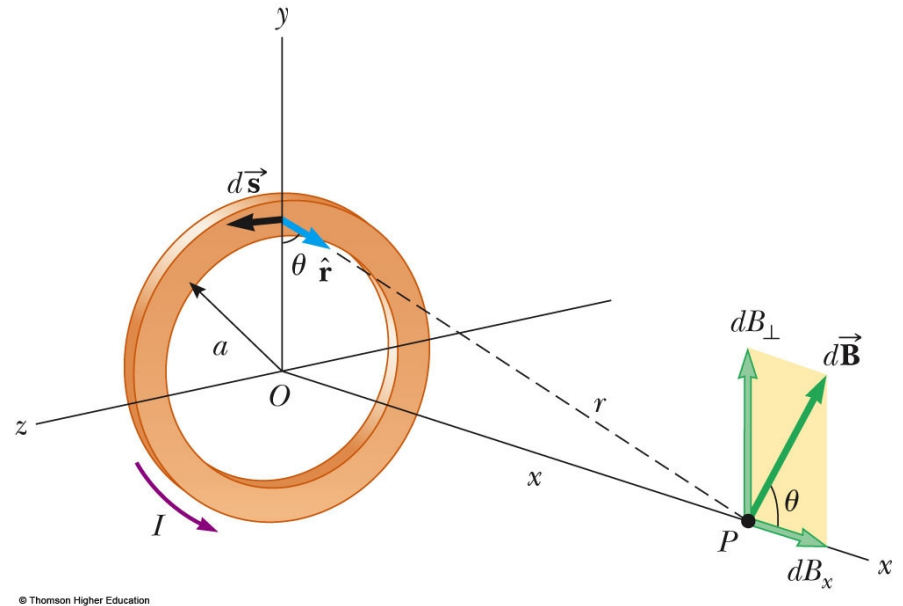
- This is the field at the *center* of the loop

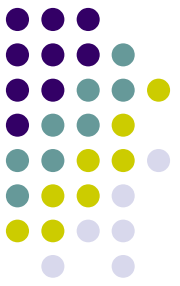


# B for a Circular Current Loop

- The loop has a radius of  $R$  and carries a steady current of  $I$
- Find the field at point  $P$

$$B_x = \frac{\mu_0 I a^2}{2(a^2 + x^2)^{3/2}}$$





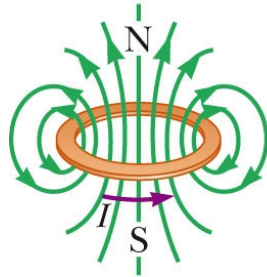
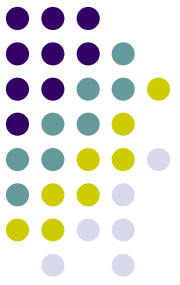
# Comparison of Loops

- Consider the field at the center of the current loop
- At this special point,  $x = 0$
- Then,

$$B_x = \frac{\mu_0 I a^2}{2(a^2 + x^2)^{3/2}} = \frac{\mu_0 I}{2a}$$

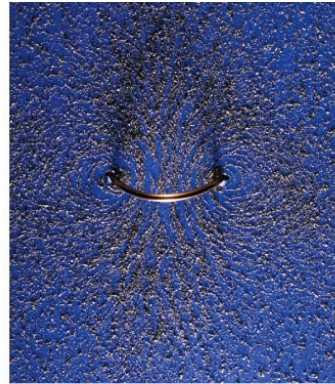
- This is exactly the same result as from the curved wire

# Magnetic Field Lines for a Loop

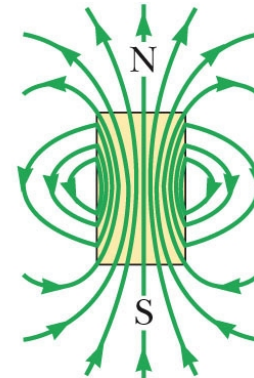


(a)

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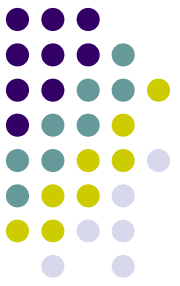
(b)



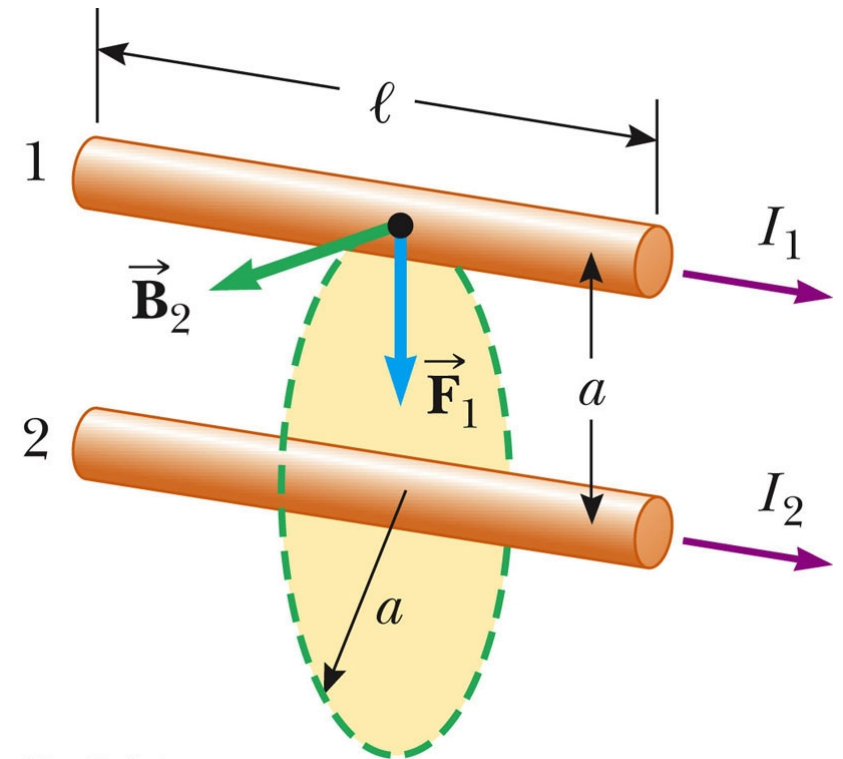
(c)

- Figure (a) shows the magnetic field lines surrounding a current loop
- Figure (b) shows the field lines in the iron filings
- Figure (c) compares the field lines to that of a bar magnet

# Magnetic Force Between Two Parallel Conductors



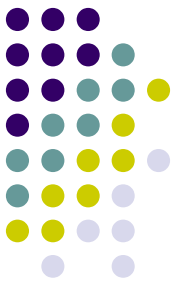
- Two parallel wires each carry a steady current
- The field  $\vec{B}_2$  due to the current in wire 2 exerts a force on wire 1 of  $F_1 = I_1 \ell B_2$



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PLAY  
ACTIVE FIGURE

# Magnetic Force Between Two Parallel Conductors, cont.

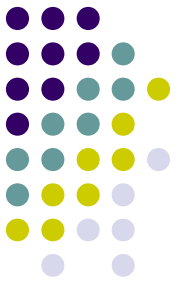


- Substituting  $\vec{B}_2 = \frac{\mu_0 I_2}{2 \pi a}$  gives

$$F_1 = \frac{\mu_0 I_1 I_2 l}{2 \pi a}$$

- RHR #1 shows us  $F_1$  is down and  $F_2$  is up.
- Parallel conductors carrying currents in the same direction attract each other
- Parallel conductors carrying current in opposite directions repel each other

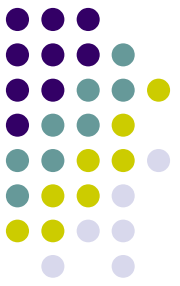
# Magnetic Force Between Two Parallel Conductors, final



- The result is often expressed as the magnetic force between the two wires,  $F_B$
- This can also be given as the force per unit length:

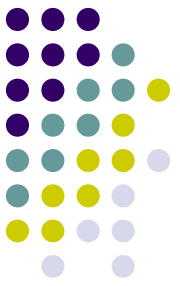
$$\frac{F_B}{l} = \frac{\mu_0 I_1 I_2}{2 \pi a}$$





# Definition of the Ampere

- The force between two parallel wires can be used to define the ampere
- When the magnitude of the force per unit length between two long, parallel wires that carry identical currents and are separated by 1 m is  $2 \times 10^{-7}$  N/m, the current in each wire is defined to be 1 A



# Definition of the Coulomb

- The SI unit of charge, the coulomb, is defined in terms of the ampere
- When a conductor carries a steady current of 1 A, the quantity of charge that flows through a cross section of the conductor in 1 s is 1 C

# Andre-Marie Ampère



- 1775 – 1836
- French physicist
- Credited with the discovery of electromagnetism
  - The relationship between electric current and magnetic fields
- Also worked in mathematics

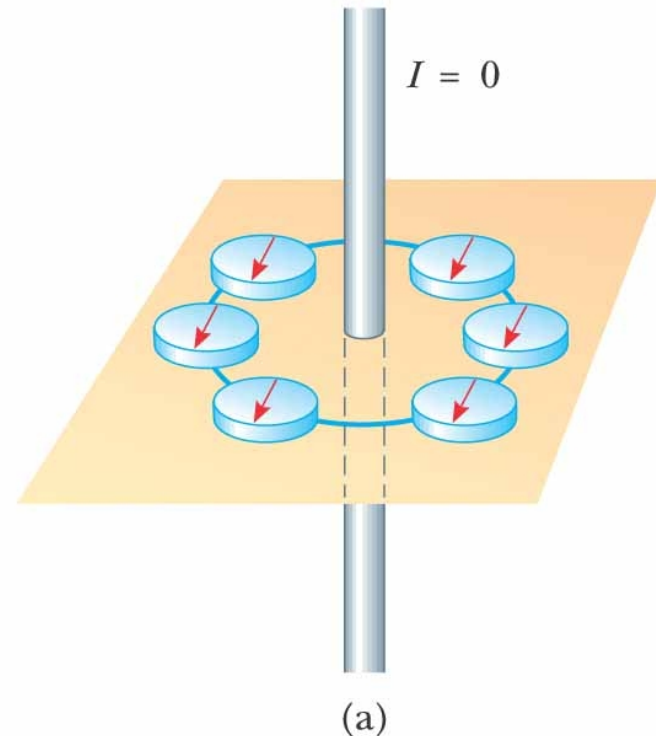


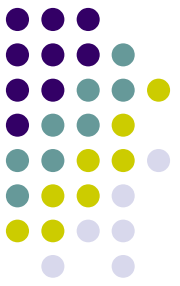
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# Magnetic Field of a Wire

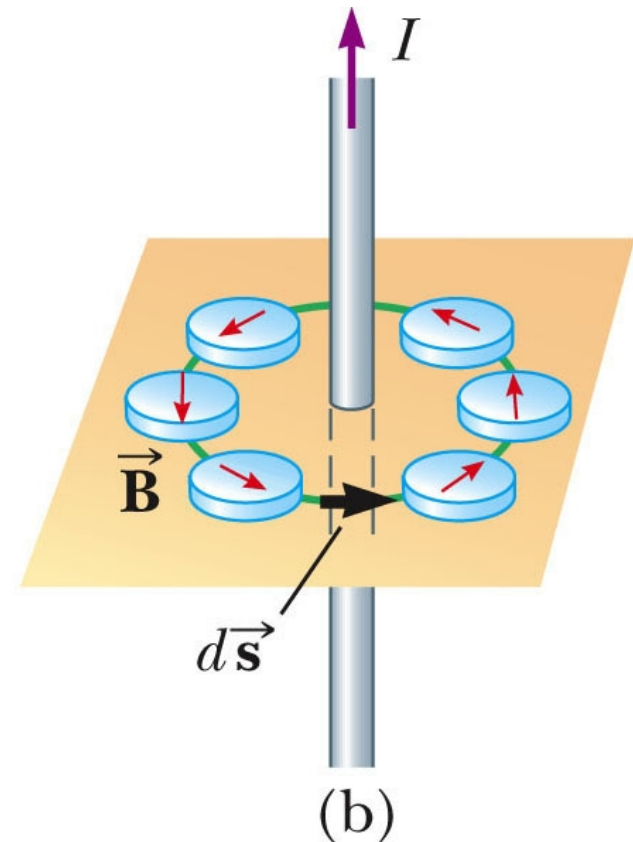
- A compass can be used to detect the magnetic field
- When there is no current in the wire, there is no field due to the current
- The compass needles all point toward the Earth's north pole
  - Due to the Earth's magnetic field





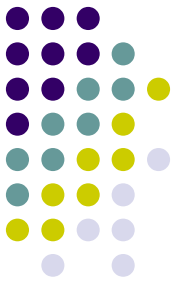
# Magnetic Field of a Wire, 2

- Here the wire carries a strong current
- The compass needles deflect in a direction tangent to the circle
- This shows the direction of the magnetic field produced by the wire
- Use the active figure to vary the current



PLAY  
ACTIVE FIGURE

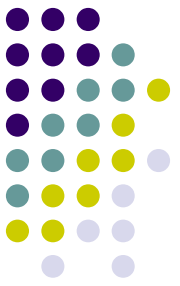
# Magnetic Field of a Wire, 3



- The circular magnetic field around the wire is shown by the iron filings



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# Ampere's Law

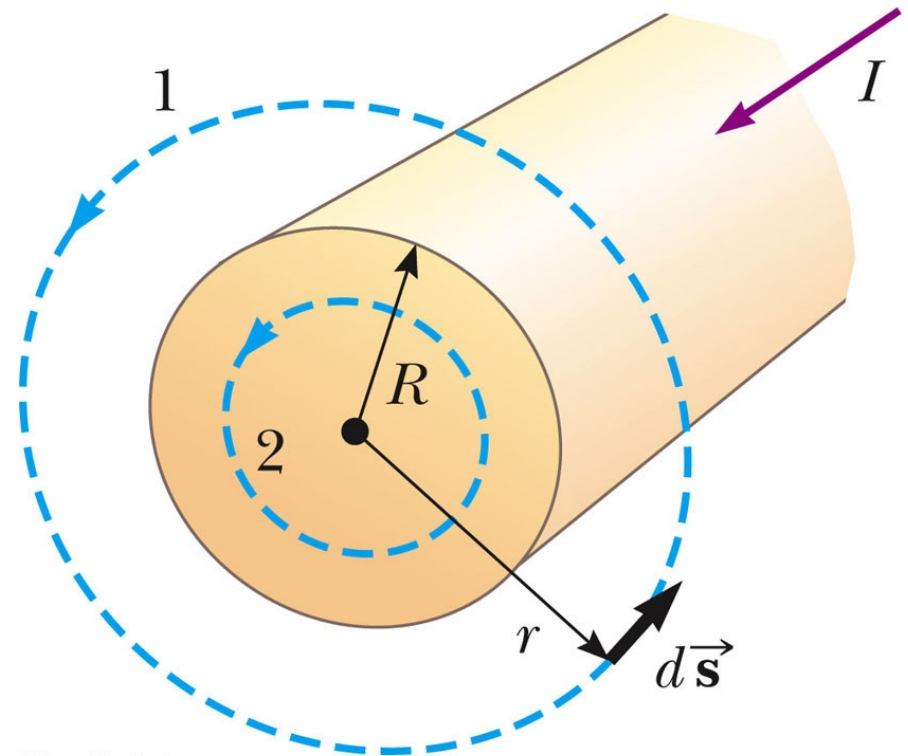
- The product of  $\vec{B} \cdot d\vec{s}$  can be evaluated for small length elements  $d\vec{s}$  on the circular path defined by the compass needles for the long straight wire
- Ampere's law states that the line integral of  $\vec{B} \cdot d\vec{s}$  around any closed path equals  $\mu_0 I_{\text{enc}}$  where  $I_{\text{enc}}$  is the total steady current passing through any surface bounded by the closed path:  
$$\vec{B} \cdot d\vec{s}$$



# Field Due to a Long Straight Wire – From Ampere’s Law



- Want to calculate the magnetic field at a distance  $r$  from the center of a wire carrying a steady current  $I$
- The current is uniformly distributed through the cross section of the wire





# Field Due to a Long Straight Wire – Results From Ampere's Law



- Outside of the wire,  $r > R$

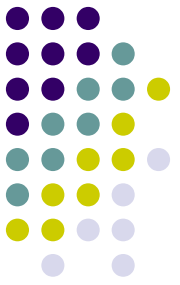
$$\int \vec{B} \cdot d\vec{s} = B(2\pi r) = \mu_0 I \qquad \int \vec{B} \cdot d\vec{s} = B(2\pi r) = \mu_0 I$$

- Inside the wire, we need  $I'$ , the current inside the amperian circle

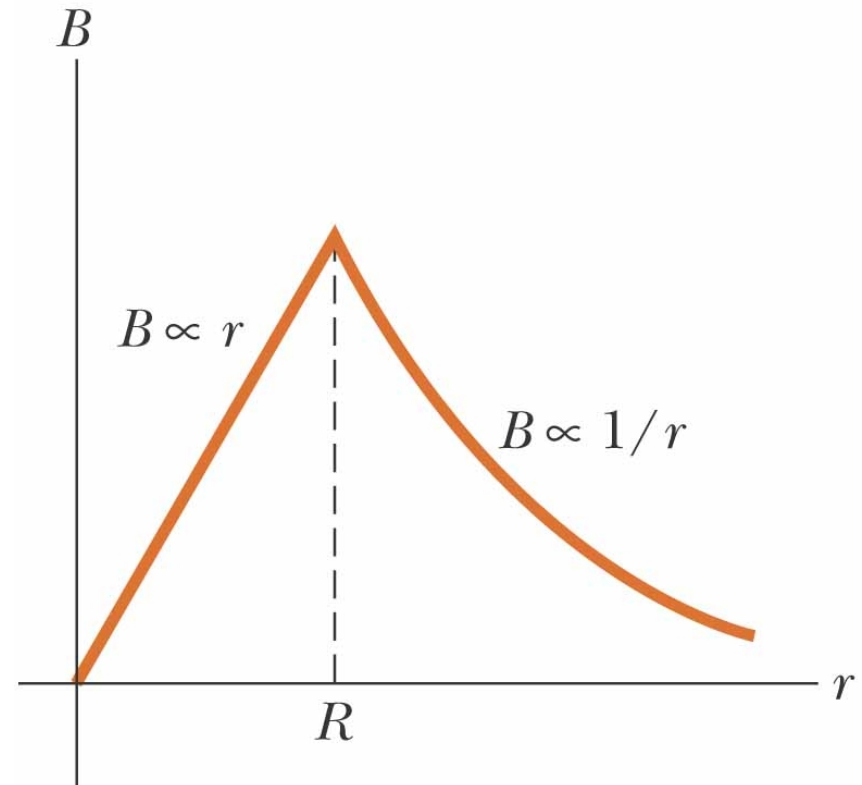
$$\int \vec{B} \cdot d\vec{s} = B(2\pi r) = \mu_0 I' \leftarrow I' = \frac{r^2}{R^2} I$$

$$B = \frac{\mu_0 I r}{2\pi R^2}$$

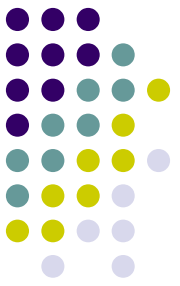
# Field Due to a Long Straight Wire – Results Summary



- The field is proportional to  $r$  inside the wire
- The field varies as  $1/r$  outside the wire
- Both equations are equal at  $r = R$



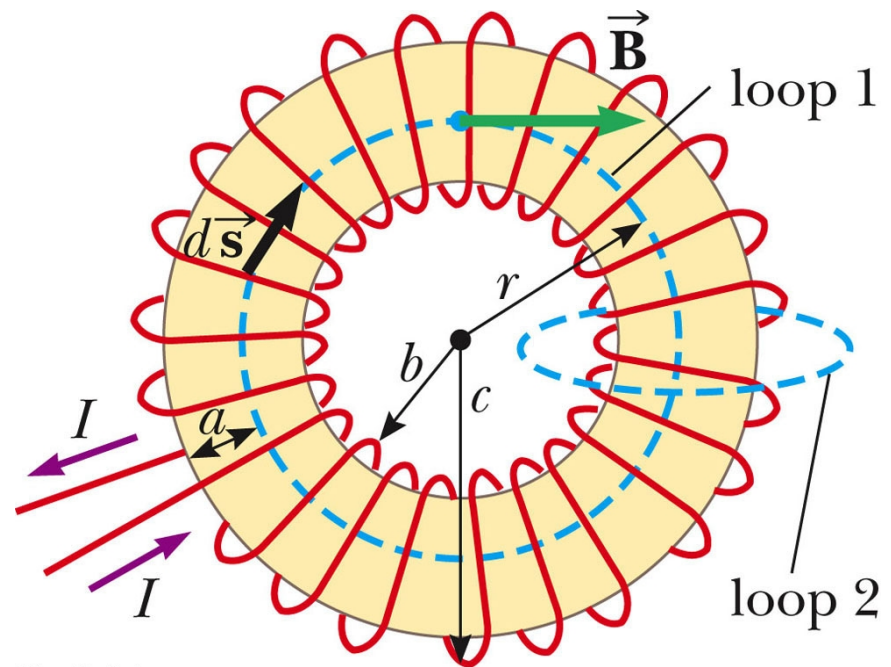
# Magnetic Field of a Toroid

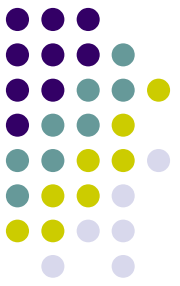


- Find the field at a point at distance  $r$  from the center of the toroid
- The toroid has  $N$  turns of wire

$$\oint \vec{B} \cdot d\vec{s} = B(2\pi r) = \mu_0 NI$$

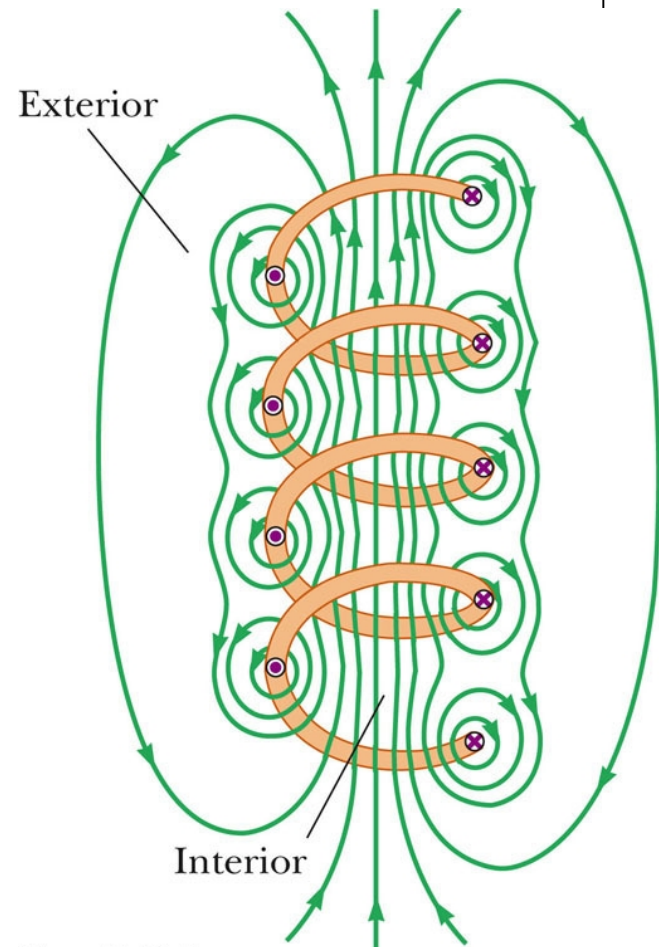
$$B = \frac{\mu_0 NI}{2\pi r}$$



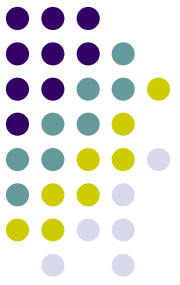


# Magnetic Field of a Solenoid

- A **solenoid** is a long wire wound in the form of a helix
- A reasonably uniform magnetic field can be produced in the space surrounded by the turns of the wire

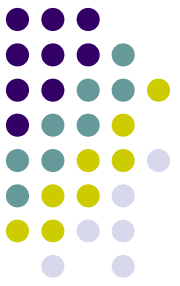


# Magnetic Field of a Solenoid, Description

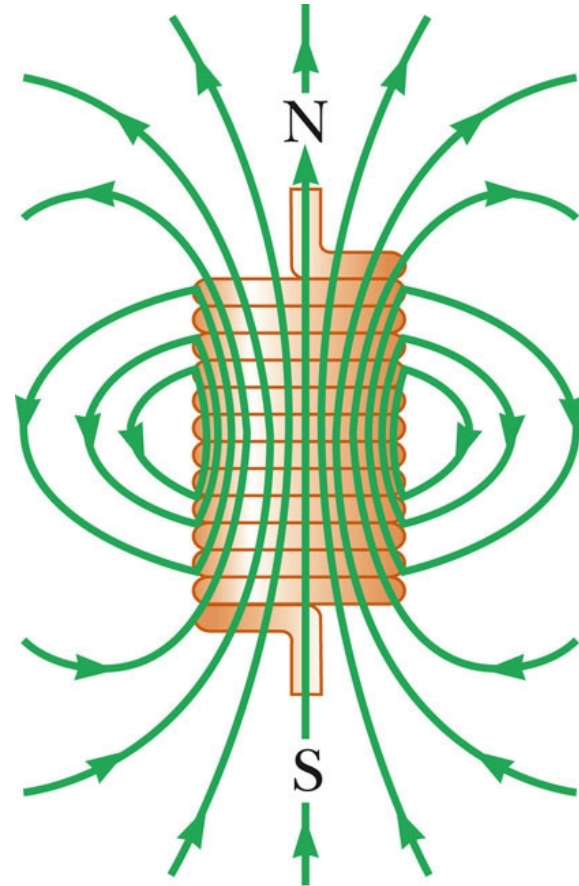


- The field lines in the interior are
  - nearly parallel to each other
  - uniformly distributed
  - close together
- This indicates the field is strong and almost uniform

# Magnetic Field of a Tightly Wound Solenoid



- The field distribution is similar to that of a bar magnet
- As the length of the solenoid increases
  - the interior field becomes more uniform
  - the exterior field becomes weaker

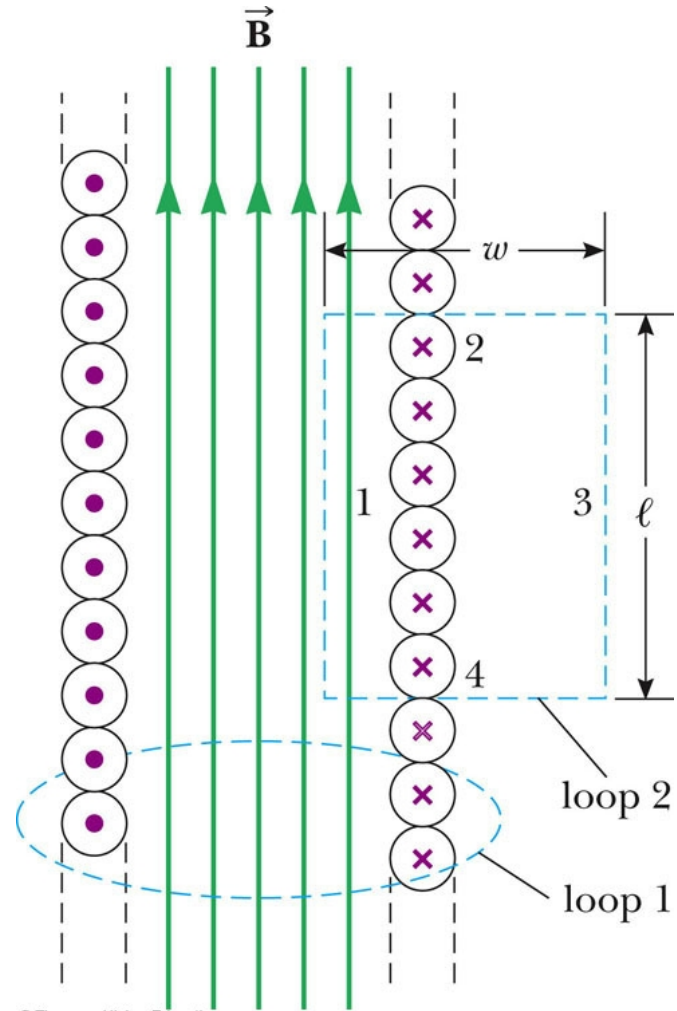


(a)

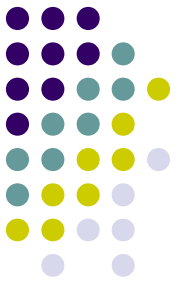
# Ideal Solenoid – Characteristics



- An *ideal solenoid* is approached when:
  - the turns are closely spaced
  - the length is much greater than the radius of the turns



# Ampere's Law Applied to a Solenoid



- Ampere's law can be used to find the interior magnetic field of the solenoid
- Consider a rectangle with side  $\ell$  parallel to the interior field and side  $w$  perpendicular to the field
  - This is loop 2 in the diagram
- Only the side of length  $\ell$  inside the solenoid contributes to the integral.
  - This is side 1 in the diagram



# Ampere's Law Applied to a Solenoid, cont.



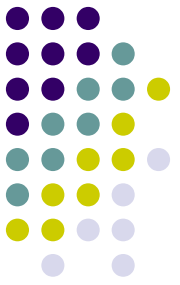
- Applying Ampere's Law gives

$$\oint \vec{B} \cdot d\vec{s} = \int_{path1} B \cdot d\vec{s} = B \int_{path1} ds = Bl$$

- The total current through the rectangular path equals the current through each turn multiplied by the number of turns

$$Bl = \mu_0 NI \dots B = \frac{\mu_0 NI}{l}$$

# Magnetic Field of a Solenoid, final



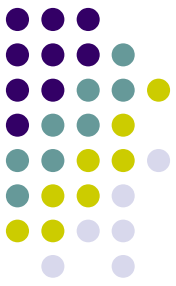
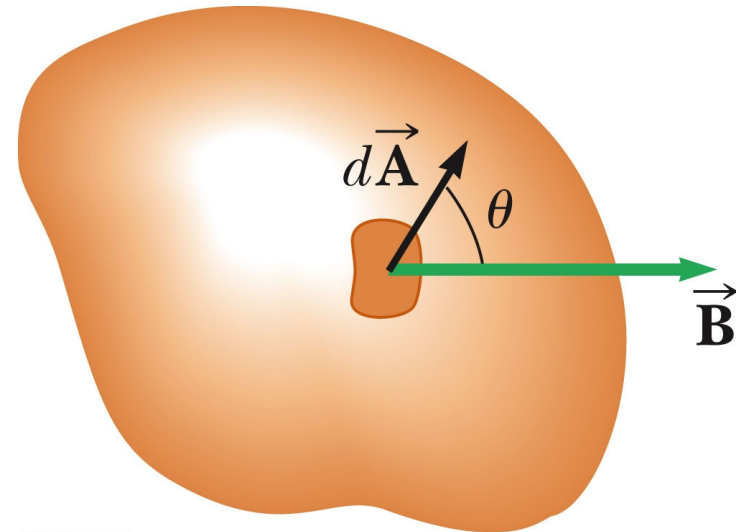
- Solving Ampere's law for the magnetic field is

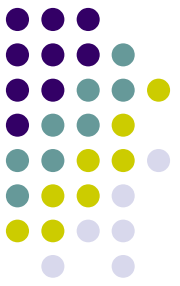
$$B = \frac{\mu_0 NI}{l} = \mu_0 nI$$

- $n = N / \ell$  is the number of turns per unit length
- This is most accurate at points near the center of a real solenoid

# Magnetic Flux

- The magnetic flux associated with a magnetic field is defined in a way similar to electric flux
- Consider an area element  $dA$  on an arbitrarily shaped surface

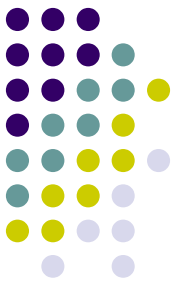




# Magnetic Flux, cont.

- The magnetic field in this element is  $\vec{B}$
- $d\vec{A}$  is a vector that is perpendicular to the surface
- $d\vec{A}$  has a magnitude equal to the area  $dA$
- The magnetic flux  $\Phi_B$  is
$$\Phi_B = \int \vec{B} \cdot d\vec{A}$$
- The unit of magnetic flux is  $\text{T} \cdot \text{m}^2 = \text{Wb}$ 
  - $\text{Wb}$  is a *weber*

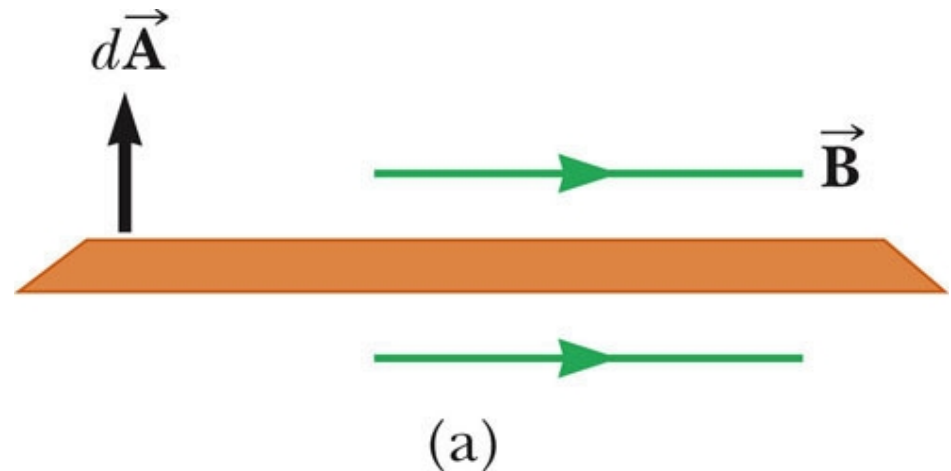
# Magnetic Flux Through a Plane, 1



- A special case is when a plane of area  $A$  makes an angle  $\theta$  with  $d\vec{A}$

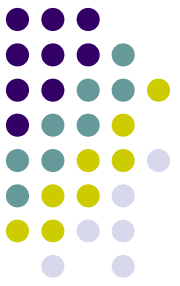
The magnetic flux is  $\Phi_B = BA \cos \theta$

- In this case, the field is parallel to the plane and  $\Phi_B = 0$



PLAY  
ACTIVE FIGURE

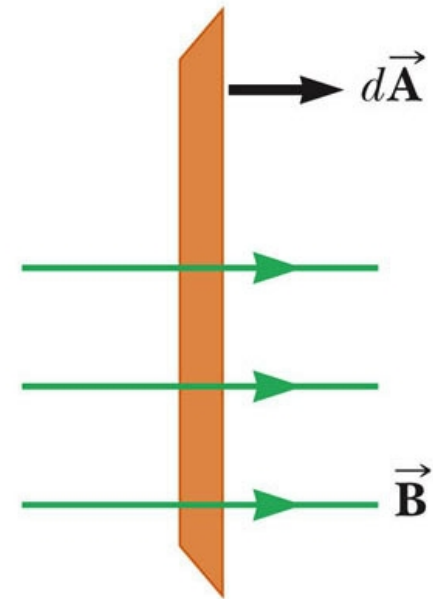
# Magnetic Flux Through A Plane, 2



- The magnetic flux is  $\Phi_B = BA \cos \theta$
- In this case, the field is perpendicular to the plane and

$$\Phi_B = BA$$

- This will be the maximum value of the flux
- Use the active figure to investigate different angles



(b)

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**PLAY  
ACTIVE FIGURE**



# Gauss' Law in Magnetism

- Magnetic fields do not begin or end at any point
  - The number of lines entering a surface equals the number of lines leaving the surface
- **Gauss' law in magnetism** says the magnetic flux through any closed surface is always zero:

$$\oint \vec{B} \cdot d\vec{A} = 0$$